

ALTERNATIVE PROOFS OF LINEAR RESPONSE FOR PIECEWISE EXPANDING UNIMODAL MAPS

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ABSTRACT. We give two new proofs that the SRB measure $t \mapsto \mu_t$ of a C^2 path f_t of unimodal piecewise expanding C^3 maps is differentiable at 0 if f_t is tangent to the topological class of f_0 . The arguments are more conceptual than the one in [4], but require proving Hölder continuity of the infinitesimal conjugacy α (a new result, of independent interest) and using spaces of bounded p -variation. The first new proof gives differentiability of higher order of $\int \psi d\mu_t$ if f_t is smooth enough and stays in the topological class of f_0 and if ψ smooth enough (a new result). In addition, this proof does not require any information on the decomposition of the SRB measure into regular and singular terms, making it potentially amenable to extensions to higher dimensions. The second new proof allows us to recover the linear response formula (i.e., the formula for the derivative at 0) obtained in [4], and gives additional information on this formula.

1. INTRODUCTION

Many chaotic dynamical systems $f : M \rightarrow M$ on a Riemannian manifold M admit an SRB measure μ (see e.g. [18]) which describes the statistical properties of a “large” set of initial conditions in the sense of Lebesgue measure. (In dimension one, an SRB measure is simply an absolutely continuous ergodic invariant probability measure $\mu_t = \rho_t dx$ with a positive Lyapunov exponent.) It is of interest (in particular in view of applications to statistical mechanics, see e.g. [14, 15]) to study the smoothness of $t \mapsto \mu_t$, when f_t is a smooth family of dynamical systems, each having an SRB measure μ_t . If $t \mapsto \mu_t$ is differentiable, one says that *linear response* holds. Ruelle [14] obtained not only differentiability, but also a formula for the derivative (the *linear response formula*), in the case of smooth uniformly hyperbolic dynamical systems.

In [4], we proved that the SRB measure $t \mapsto \mu_t$ of a C^2 family of piecewise C^3 and piecewise expanding unimodal maps f_t , with f_0 mixing (see §2.1 for formal definitions), is differentiable at $t = 0$ (as a Radon measure) if and only if f_t is tangent to the topological class of f_0 at $t = 0$. (Keller [10] proved a long time ago that ρ_t has a $|t| \ln |t|$ modulus of continuity, as an element of $L^1(dx)$, and examples in [1, 4] show that this can be attained for non tangential families.) We also obtained in [4] a linear response formula analogous to the one in [14] (we used the resummation introduced in [1]).

More recently, differentiability of the SRB measure (in the weak sense, that is, as an appropriate distribution) was obtained [16], [6] for smooth families f_t of analytic

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and nonuniformly expanding unimodal maps which stay in the topological class of f_0 . The cases of families of smooth nonuniformly expanding interval maps only tangent to the topological class (where Whitney differentiability is expected on suitable subsets of parameters), as well as higher-dimensional dynamical systems such as piecewise expanding/hyperbolic maps or Hénon-like maps, are still open, and much more difficult, see [2] for a discussion. In particular, the arguments in [4] and [16] used detailed information about the structure of the SRB measure, decomposing it into a regular and a singular term. This type of information may be far less accessible in higher dimensions.

In this article, we give two new proofs of the fact (Theorem 5.1 in [4]) that the SRB measure of a C^2 family of piecewise C^3 and piecewise expanding unimodal maps f_t , with f_0 mixing, is differentiable at $t = 0$ if f_t is tangent to the topological class of f_0 at $t = 0$.

Section 3 contains our first new proof (see Corollary 3.2), more precisely, we obtain differentiability of $t \mapsto \int \psi d\mu_t$ for $\psi \in C^{1+\text{Lip}}$ if f_t is a C^2 family of piecewise expanding C^3 unimodal maps tangent to the topological class of a mixing map f_0 . The argument is based on thermodynamic formalism, using potentials $(s, t) \mapsto s(\psi \circ h_t) - \log |f'_t \circ h_t|$ (where h_t conjugates \tilde{f}_t with f_0 and $|\tilde{f}_t - f_t| = O(t^2)$) and does not require any knowledge about the structure of μ_t . It may therefore be useful in more difficult situations (such as Hénon maps, see [2]). It requires proving Hölder differentiability of the *infinitesimal conjugacy* α , a new result (Proposition 2.3), of independent interest. Also, this new proof gives that $t \mapsto \int \psi d\mu_t$ is a C^j function, if $\psi \in C^{j+\text{Lip}}$ and f_t is a C^{j+1} family of piecewise expanding C^{j+2} maps in the topological class of f_0 , for any $j \geq 1$ (this is a new result, Theorem 3.1). Note also that we do not require the assumption from [4] that there is a function X so that $\partial f_t|_{t=0} = X \circ f_0$.

The first new proof requires $\psi \in C^{1+\text{Lip}}$ (instead of $\psi \in C^0$ as in [4]) and does not furnish the linear response formula. Section 4 contains our second new proof (Theorem 4.1), which uses spectral perturbation theory for transfer operators associated to the dynamics f_0 and the weight $1/|f'_t \circ h_t|$. This other proof gives differentiability of $\int \psi d\mu_t$ for $\psi \in C^0$ and, using the assumption that $\partial f_t|_{t=0} = X \circ f_0$, allows us to recover the linear response formula from [4]. (This second proof also uses the Hölder regularity of α from Proposition 2.3.) Note however that this second proof requires information on the structure of μ_t from [1, Prop. 3.3].

Putting together Theorem 3.1 and 4.1 we get the following additional result (Corollary 4.4): If f_t is a C^2 family of piecewise expanding C^3 unimodal maps in the topological class of f_0 , then $t \mapsto \mu_t$ is C^1 from a neighbourhood of zero to Radon measures.

We emphasize that neither new proof gives that the condition to be tangent to the topological class is necessary, contrary to the argument in [4] (see Theorem 7.1 there). The proofs here are a bit shorter than the one given in [4], although the present account requires some results from our previous papers (such as [1, Prop. 3.3] [5, Prop. 3.2, Theorem 2] [4, Prop. 2.4, Lemma 2.6, Prop. 3.3]).

2. DEFINITIONS AND NOTATIONS – HÖLDER SMOOTHNESS OF THE INFINITESIMAL CONJUGACY

2.1. Formal definitions. Denote $I = [-1, 1]$. For an integer $k \geq 1$, we define \mathcal{B}^k to be the linear space of continuous functions $f: I \rightarrow \mathbb{R}$ such that f is C^k on the intervals $[-1, 0]$ and $[0, 1]$. Then \mathcal{B}^k is a Banach space for the norm $\max\{|f|_{C^k([-1, 0])}, |f|_{C^k([0, 1])}\}$. For an integer $k \geq 1$, we define the set \mathcal{U}^k of *piecewise expanding C^k unimodal maps* to be the set of $f \in \mathcal{B}^k$ such that ¹ $f(-1) = f(1) = -1$, $\inf_{x \neq 0} |f'(x)| > 1$, and $f(0) \leq 1$ (so that $f(I) \subset I$). The point $c = 0$ is called the *critical point* of f .

A piecewise expanding C^k unimodal map f is *good* if either c is not periodic under f or $|(f^{q-1})'(f(c))| \min\{|f'_+(c)|, |f'_-(c)|\} > 2$, where $q \geq 2$ is the minimal period of c ; it is *mixing* if f is topologically mixing on $[c_2, c_1]$, where $c_k = f^k(c)$.

For $1 \leq j \leq k$, a C^j *family of piecewise expanding C^k unimodal maps* is a C^j map f_t from $t \in (-\epsilon, \epsilon)$ to \mathcal{U}^k for some $\epsilon > 0$. (In this paper, $k \geq 1$ is an integer and j is either an integer or $j = k - 1 + \text{Lip}$ for $k \geq 2$, the notations $\mathcal{B}^{k+\text{Lip}}$ and $\mathcal{U}^{k+\text{Lip}}$ for integers $k \geq 1$ being self-explanatory. See also Remark 2.2.)

Remark 2.1. A C^j family f_t of piecewise expanding C^k unimodal maps is a $C^{j,k}$ perturbation of f_0 in the sense of [4] if $j = k \geq 2$.

Remark 2.2. Considering $\mathcal{B}^{k+\beta}$ and $\mathcal{U}^{k+\beta}$ for $k \geq 1$ integer and a Hölder exponent $0 < \beta < 1$ will perhaps allow to avoid the loss of regularity from C^{k+1} to $C^{k-1+\text{Lip}}$ e.g. in [4, Prop. 2.4] (this question was asked by J.-C. Yoccoz). However, since the spectral result of Wong [17] only holds on the space BV_p of functions of bounded p -variation if $1 \leq p < p_0$, for some $p_0 > 1$ depending on the dynamics, it may be necessary in this case to replace BV_p by spaces of generalised p -variation, as introduced by Keller [11]. (See also Remark 2.5.)

Assume that f_t is a C^j family of piecewise expanding C^k unimodal maps for $k \geq j > 1$. By classical results of Lasota–Yorke, each f_t has a unique absolutely continuous invariant probability measure $\mu_t = \rho_t dx$. This measure is ergodic and it is called the *SRB measure* of f_t . If f_t is mixing, then μ_t is mixing. If f_0 is good and mixing, then f_t is mixing for all small enough t (see [10] and references therein).

We say that a C^j family f_t of piecewise expanding C^k unimodal maps is in the *topological class* of f_0 if there exist homeomorphisms $h_t: I \rightarrow I$ such that

$$(1) \quad h_0 = \text{id} \text{ and } h_t \circ f_0 = f_t \circ h_t, \quad \forall |t| < \epsilon.$$

We say that a C^j family f_t of piecewise expanding C^k unimodal maps ($k \geq j \geq 2$) is *tangent to the topological class* of f_0 if there exists a C^j family \tilde{f}_t of piecewise expanding C^k unimodal maps in the topological class of f_0 so that $\tilde{f}_t = f_0$ and $\partial_t f_t|_{t=0} = \partial_t \tilde{f}_t|_{t=0}$. (Note that there is a typographical mistake in [4, p. 682, line 6], where “ $C^{2,2}$ perturbation” should be replaced by “ $C^{r_0, r}$ perturbation.”)

We say that a bounded function $v: I \rightarrow \mathbb{R}$ is *horizontal* for f , if $v(-1) = v(1) = 0$, and setting $M_f = q$ if c is periodic of minimal period q , and $M_f = +\infty$ otherwise,

$$(2) \quad J(f, v) = \sum_{j=0}^{M_f-1} \frac{v(c_j)}{(f^j)'(c_1)} = 0.$$

¹A prime denotes differentiation with respect to $x \in I$, a priori in the sense of distributions.

In [4, Cor. 2.6] we proved that if f_t is a C^2 family of piecewise expanding C^2 unimodal maps tangent to the topological class of f_0 , then $v = \partial f_t|_{t=0}$ is horizontal for f_0 . By [4, Theorem 2], if f_t is a C^2 family of piecewise expanding C^2 unimodal maps with f_0 good and $v = \partial f_t|_{t=0}$ is C^2 and horizontal for f_0 , then there exists a $C^{1+\text{Lip}}$ family \tilde{f}_t of piecewise expanding C^2 unimodal maps in the topological class of f_0 so that $\tilde{f}_t = f_0$ and $\partial_t \tilde{f}_t|_{t=0} = \partial_t f_t|_{t=0}$.

We proved in [4, Lemma 2.2] that if $v : I \rightarrow \mathbb{R}$ is bounded then the twisted cohomological equation (TCE)

$$(3) \quad v(x) = \alpha(f(x)) - f'(x)\alpha(x), \quad \forall x \in I, \quad x \neq c,$$

admits a unique bounded solution α satisfying $\alpha(c) = 0$. This solution is obtained as follows: If c is not in the forward orbit of x , set $M(x) = \infty$ and otherwise let $M(x)$ be the smallest integer $j \geq 0$ satisfying $f^j(x) = c$, then put

$$(4) \quad \alpha(x) = \sum_{i=0}^{M(x)-1} \frac{v(f^i(x))}{(f^i)'(x)}.$$

The function α is called the *infinitesimal conjugacy*.

If $u : I \rightarrow \mathbb{R}$ is Hölder we denote its Hölder norm by $|u|_\beta$. Slightly abusing notation, we shall sometimes write $\partial_t f_t$ for $\partial_s f_s|_{s=t}$, and similarly for other functions depending on t .

2.2. Hölder smoothness of the infinitesimal conjugacy α . A new result that we shall require throughout (see Lemmas 3.3 and 4.2) is:

Proposition 2.3 (Smoothness of the infinitesimal conjugacy). *Let $f \in \mathcal{U}^2$ be such that c is not periodic. For any $\beta \in (0, 1)$ there exist $C_\beta > 0$ and \mathcal{V}_β a neighbourhood of f in \mathcal{U}^2 so that, for any $g \in \mathcal{V}_\beta$ and every β -Hölder $v : I \rightarrow \mathbb{R}$, with $v(-1) = v(1) = 0$ and $J(g, v) = 0$, the unique bounded function α (4) satisfying $\alpha(c) = 0$ and $v(x) = \alpha(g(x)) - g'(x)\alpha(x)$ for all $x \neq c$ is β -Hölder, with*

$$|\alpha|_\beta \leq C_\beta |v|_\beta.$$

If the critical point of $f \in \mathcal{U}^2$ is periodic, the statement holds up to taking (for appropriate $\xi(\beta) > 0$)

$$\mathcal{V}_\beta = \{g \mid \|g - f\|_{\mathcal{B}^2} < \xi(\beta), \exists \text{ homeomorphism } h : I \rightarrow I \text{ s.t. } g \circ h = h \circ f\}.$$

In particular, if $f \in \mathcal{U}^2$ and $J(f, v) = 0$ for some Lipschitz v with $v(-1) = v(1) = 0$, the function α solving (3) is β -Hölder for any $\beta < 1$.

Remark 2.4. Jérôme Buzzi [8] showed us a simple proof that if h is a homeomorphism so that $h \circ f = g \circ h$, for two piecewise expanding C^1 unimodal maps f and g , then h is β -Hölder, for any $\beta < \log(\inf |g'|/2) / \log(2 \sup |f'|)$. This fact neither implies nor is implied by Proposition 2.3.

Proof. Step I. For any $\beta < 1$, there exist a neighbourhood \mathcal{V}_β of f in \mathcal{U}^2 , $\ell \geq 1$, and $\eta > 0$ such that $\lambda = (\inf_{g \in \mathcal{V}_\beta} \inf_{x \neq c} |g'(x)|)^{-1} < 1$, and, for any $g \in \mathcal{V}_\beta$, letting $d_1 < d_2 < \dots < d_p$ be the critical points of g^ℓ , putting $d_0 = -1$, $d_{p+1} = 1$, and setting

$$(5) \quad \theta = \max_{0 \leq i \leq p} \sup_{\substack{x, y \in (d_i, d_{i+1}) \\ |x-y| < \eta}} \frac{|(g^\ell)'(x)|^\beta}{|(g^\ell)'(y)|},$$

we have $2\theta < 1$.

Put $\Delta_g = \min_{0 \leq i \leq p} \{d_{i+1} - d_i\}$. Then $\inf_{g \in \mathcal{V}_\beta} \Delta_g > 0$ if the critical point of f is not periodic. Otherwise we have $\inf_{g \in \mathcal{V}_\beta} \Delta_g > 0$, up to replacing \mathcal{V}_β by a \mathcal{B}^2 -neighbourhood of f in its topological class. In particular, we can assume that $\eta < \inf_{g \in \mathcal{V}_\beta} \Delta_g$. From now on, we fix \mathcal{V}_β , $\ell \geq 1$, and $\eta > 0$ as above.

Step II. We claim it suffices to show the lemma for $g \in \mathcal{V}_\beta$ with a periodic critical point: Indeed, if g has a nonperiodic critical point, then we consider $g_t = g + tw$ with $g_t \in \mathcal{U}^2$, $w \in \mathcal{B}^2$, $w(-1) = w(1) = 0$, and $J(g, w) \neq 0$. By [5, Corollary 4.1], there exists a sequence $t_n \rightarrow 0$ such that each $g_n = g_{t_n}$ has a periodic critical point. In particular, g_n converges to g in the \mathcal{U}^2 topology. Then, by [5, Proposition 3.2] we have $\lim_{n \rightarrow \infty} J(g_n, v) = 0$. Let w_n be a β -Hölder function, with $w_n(-1) = w_n(1) = 0$ and $|w_n|_\beta \leq 1$, such that $J(g_n, w_n) = 1$. Set

$$v_n = v - J(g_n, v)w_n.$$

Then we have $J(g_n, v_n) = 0$ and $\lim_{n \rightarrow \infty} |v_n - v|_\beta = 0$. If the proposition holds for maps in \mathcal{V}_β with a periodic turning point, the unique function α_n so that $\alpha_n(c) = 0$ and $v_n(x) = \alpha_n(g_n(x)) - g'_n(x)\alpha_n(x)$ for all $x \neq c$, satisfies $|\alpha_n|_\beta \leq C_\beta |v_n|_\beta$. We can choose a subsequence α_{n_i} converging in the sup norm to a function α . It follows from the uniform convergence of α_{n_i} that α satisfies the TCE (3) for g and v , and that $|\alpha|_\beta \leq C_\beta |v|_\beta$.

Step III. We assume from now on that $g \in \mathcal{V}_\beta$ has a periodic turning point. The proof will be via an “infinitesimal pull-back” argument.

First, since $J(g, v) = 0$, it is easy to see that there exists a β -Hölder function $\alpha_0 : I \rightarrow \mathbb{R}$ with $\alpha_0(-1) = \alpha_0(1) = \alpha_0(c) = 0$, $\alpha_0(g(c)) = v(c)$, and

$$v(x) = \alpha_0(g(x)) - g'(x)\alpha_0(x) \text{ for every } x \neq c \text{ in the (finite) forward orbit of } c.$$

Second, we define by induction continuous functions $\alpha_i : I \rightarrow \mathbb{R}$, for $i \geq 1$, such that $\alpha_i(-1) = \alpha_i(1) = \alpha_i(c) = 0$, $\alpha_i(g(c)) = v(g(c))$, that

$$(6) \quad v(x) = \alpha_i(g(x)) - g'(x)\alpha_i(x) \text{ for every } x \neq c \text{ in the (finite) forward orbit of } c,$$

and, in addition,

$$(7) \quad v(x) = \alpha_{i-1}(g(x)) - g'(x)\alpha_i(x), \forall x \neq c.$$

Indeed, suppose we have defined α_i , for $0 \leq i \leq n$. Set $\alpha_{n+1}(c) = 0$, and

$$\alpha_{n+1}(x) = \frac{\alpha_n(g(x)) - v(x)}{g'(x)}, \quad x \neq c.$$

Clearly, $\alpha_{n+1}(-1) = \alpha_{n+1}(1) = 0$, and (7) holds for $i = n+1$. Thus, since $v(x) = \alpha_n(g(x)) - g'(x)\alpha_n(x)$ for every $x \neq c$ in the forward orbit of c , we find $\alpha_n(x) = \alpha_{n+1}(x)$ for each $x \neq c$ in the forward orbit of c . Since $\alpha_n(c) = \alpha_{n+1}(c) = 0$, we conclude that (6) holds for $i = n+1$, and $\alpha_{n+1}(g(x)) = v(x)$. Last, but not least, α_{n+1} is continuous on I because $\alpha_n(g(x)) = v(x)$.

Thirdly, if x is not a critical point of g^j , we set $v_0(x) = 0$, and

$$v_j(x) = \sum_{i=0}^{j-1} (g^{j-1-i})'(g^{i+1}(x))v(g^i(x)), \quad j \geq 1.$$

Recalling the notation ℓ , $\{d_i\}$, from Step I, it is easy to see that

$$(8) \quad v_\ell(x) = \alpha_{j\ell}(g^\ell(x)) - (g^\ell)'(x)\alpha_{(j+1)\ell}(x), \quad \forall x \notin \{d_1, \dots, d_p\}, \forall j \geq 0.$$

For $\ell \geq 2$ the function v_ℓ may have jump discontinuities at the critical points d_i of g^ℓ , but it is β -Hölder in the connected components of $I \setminus \{d_1, \dots, d_p\}$.

Finally, we shall use the iterated twisted cohomological equation (8) to show that there exists $C_\beta < \infty$ so that, for all $g \in \mathcal{V}_\beta$ with a periodic turning point and all β -Hölder v with $J(g, v) = 0$ (and $v(-1) = v(1) = 0$), there exists $j_0 \geq 0$ so that $|\alpha_{j\ell}|_\beta \leq C_\beta |v|_\beta$ for all $j \geq j_0$. In view of this, for $j \geq 0$, set $K_j^0 = \sup_I |\alpha_{j\ell}|$, $L^0 = \sup_I |v_\ell|$,

$$K_j^\beta = \sup_{x \neq y} \frac{|\alpha_{j\ell}(x) - \alpha_{j\ell}(y)|}{|x - y|^\beta}, \quad \widehat{K}_j^\beta = \max_{0 \leq i \leq p} \sup_{\substack{x \neq y, |x-y| < \eta \\ x, y \in (d_i, d_{i+1})}} \frac{|\alpha_{j\ell}(x) - \alpha_{j\ell}(y)|}{|x - y|^\beta},$$

and

$$L^\beta = \max_{0 \leq i \leq p} \sup_{\substack{x \neq y \\ x, y \in (d_i, d_{i+1})}} \frac{|v_\ell(x) - v_\ell(y)|}{|x - y|^\beta}, \quad D = \max_{0 \leq i \leq p} \sup_{x, y \in (d_i, d_{i+1})} \frac{|(g^\ell)''(x)|}{|(g^\ell)'(y)|^2}.$$

Clearly, $\max(L^0, L^\beta) \leq \widetilde{C}_\beta |v|_\beta$ for all g and v under consideration, and we have

$$(9) \quad \frac{|\alpha_{(j+1)\ell}(x) - \alpha_{(j+1)\ell}(y)|}{|x - y|^\beta} \leq 2\eta^{-\beta} K_{j+1}^0 \quad \text{if } |x - y| \geq \eta, \forall j \geq 0.$$

Therefore, recalling the definition of λ and (5) from Step I, it suffices to show that

$$(10) \quad K_{j+1}^0 \leq \lambda^\ell (K_j^0 + L^0) \text{ and } \widehat{K}_{j+1}^\beta \leq (L^0 + K_j^0)D + \lambda^\ell L^\beta + \theta K_j^\beta, \forall j \geq 0.$$

Indeed, continuity of $\alpha_{(j+1)\ell}$ together with (9) (recall also that $\eta < \inf_g \Delta_g$, so that if $|x - y| < \eta$ then $[x, y]$ contains at most one point d_i) imply

$$K_{j+1}^\beta \leq \max(2\eta^{-\beta} K_{j+1}^0, 2\widehat{K}_{j+1}^\beta).$$

The above bound together with (10) yield $E_\beta < \infty$ so that, for all $g \in \mathcal{V}_\beta$ with a periodic turning point, and all β -Hölder v with $J(g, v) = 0$, $v(-1) = v(1) = 0$, there exists $j_0 \geq 0$ so that

$$K_{j+1}^\beta \leq E_\beta |v|_\beta + 2\theta K_j^\beta, \forall j \geq j_0.$$

Since $2\theta < 1$, we conclude by a geometric series, taking larger j_0 if necessary.

It remains to show (10). We concentrate on the second bound (the first is easier and left to the reader). Let $x, y \in (d_i, d_{i+1})$ satisfy $|x - y| < \eta$. Then (8) implies (since $g \in \mathcal{U}^2$, the function $(g^\ell)'$ is C^1 in the intervals of monotonicity of g^ℓ)

$$\begin{aligned} |\alpha_{(j+1)\ell}(x) - \alpha_{(j+1)\ell}(y)| &\leq |v_\ell(x) + \alpha_{j\ell}(g^\ell(x))| \left| \frac{1}{(g^\ell)'(x)} - \frac{1}{(g^\ell)'(y)} \right| \\ &\quad + \frac{|v_\ell(x) - v_\ell(y)| + |\alpha_{j\ell}(g^\ell(x)) - \alpha_{j\ell}(g^\ell(y))|}{|(g^\ell)'(y)|} \\ &\leq (L^0 + K_j^0) \frac{\sup_{[x, y]} |(g^\ell)''|}{\inf_{[x, y]} |(g^\ell)'|^2} |x - y| + \frac{L^\beta + \sup_{[x, y]} |(g^\ell)'|^\beta K_j^\beta}{\inf_{[x, y]} |(g^\ell)'|} |x - y|^\beta \\ &\leq (L^0 + K_j^0)D |x - y| + (\lambda^\ell L^\beta + \theta K_j^\beta) |x - y|^\beta. \end{aligned}$$

Step IV. Defining $\tilde{\alpha}_n = \frac{1}{n} \sum_{j=0}^{n-1} \alpha_{j\ell}$, we can choose a subsequence $\tilde{\alpha}_{n_i}$ converging uniformly on I to a function $\tilde{\alpha}$ satisfying $|\tilde{\alpha}|_\beta \leq C_\beta |v|_\beta$. By (8),

$$(11) \quad v_\ell(x) = \tilde{\alpha}(g^\ell(x)) - (g^\ell)'(x) \tilde{\alpha}(x), \quad \forall x \notin \{d_1, \dots, d_p\}.$$

Let $\alpha: I \rightarrow \mathbb{R}$ be the unique bounded solution vanishing at c to the TCE (3) for g and v , as in (4). Then

$$(12) \quad v_\ell(x) = \alpha(g^\ell(x)) - (g^\ell)'(x)\alpha(x), \quad \forall x \notin \{d_1, \dots, d_p\}.$$

Since $\tilde{\alpha}$ is continuous (11) and (12), imply $\alpha = \tilde{\alpha}$. We proved $|\alpha|_\beta \leq C_\beta |v|_\beta$, for all $g \in \mathcal{V}_\beta$ with a periodic turning point, and thus the proposition. \square

2.3. Banach spaces of bounded variation. We shall consider the Banach space of functions of bounded variation

$$BV = \{\varphi: \mathbb{R} \rightarrow \mathbb{C} \mid \text{var}(\varphi) < \infty, \text{supp}(\varphi) \subset I\} / \sim,$$

endowed with the norm $\|\varphi\|_{BV} = \inf_{\psi \sim \varphi} \text{var}(\psi)$, where var denotes total variation, and $\varphi_1 \sim \varphi_2$ if the bounded functions φ_1, φ_2 differ on an at most countable set. In addition, for $1 \leq p < \infty$ we shall work with the Banach space of functions of bounded p variation (used in interval dynamics by Wong [17])

$$BV_p = \{\varphi: \mathbb{R} \rightarrow \mathbb{C} \mid \text{var}_p(\varphi) < \infty, \text{supp}(\varphi) \subset I\} / \sim,$$

where

$$\text{var}_p(\varphi) = \sup_{x_1 < x_2 < \dots < x_n} \left(\sum_{i=1}^n |\varphi(x_{i+1}) - \varphi(x_i)|^p \right)^{1/p},$$

the supremum ranging over all ordered finite subsets of \mathbb{R} . Note that $\text{var}_1 = \text{var}$ and $BV = BV_1$. Wong [17] does not quotient by the equivalence relation $\varphi_1 \sim \varphi_2$, but his results remain unchanged if we consider elements in BV_p modulo \sim (a function in BV_p is continuous except on an at most countable set, see also [11, Lemma 1.4.a, Lemma 2.7] and [7]). Note that for each $p \geq 1$ there is $C \geq 1$ so that $|\varphi|_\infty \leq C \|\varphi\|_{BV_p}$ for all φ , and if φ is $1/p$ -Hölder, then $\|\varphi\|_{BV_p} \leq |\varphi|_{1/p}$. Also,

$$(13) \quad \|\varphi_1 \varphi_2\|_{BV_p} \leq 2 \|\varphi_1\|_{BV_p} \|\varphi_2\|_{BV_p}, \quad \forall p \geq 1,$$

and $\|\varphi \circ h\|_{BV_p} = \|\varphi\|_{BV_p}$ for any homeomorphism $h: I \rightarrow I$ and all $p \geq 1$.

Remark 2.5. The reason we consider spaces BV_p for $p \neq 1$ is because we are concerned with differentiability in the t -parameter and we will have to deal with derivatives $\partial_t(\psi \circ h_t)|_{t=0} = \psi' \cdot \partial_t h_t|_{t=0}$ or $\partial_t(f'_t \circ h_t)|_{t=0} = f''_0 \partial_t h_t|_{t=0} + v'$, where $v' = \partial_t f'_t|_{t=0}$ is C^1 , but $\partial_t h_t|_{t=0}$ does not belong to BV in general. We shall see, however, that Proposition 2.3 implies that $\alpha = \partial_t h_t|_{t=0}$ lies in BV_p for all $p > 1$.

3. WEAK DIFFERENTIABILITY OF THE SRB VIA THE PRESSURE

The main result of this section (Theorem 3.1) says that for any $j \geq 1$, if f_t is a C^{j+1} family of piecewise expanding C^{j+2} unimodal maps in the topological class of f_0 , then $R(t) = \int \psi d\mu_t$ is C^j if ψ is $C^{j+\text{Lip}}$. Even if $j = 1$, this is a new result (Theorem 5.1 in [4] only gives differentiability at $t = 0$). The argument is based on the topological pressure of the potential $(s, t) \mapsto -\log |f'_t \circ h_t| + s(\psi \circ h_t)$ for the map f_0 . It is simple, but does not give the formula for $\partial_t R(t)|_{t=0}$ (or higher order derivatives). Using the linear response formula from [4, Theorem 5.1] or Theorem 4.1 below, Theorem 3.1 will imply Corollary 4.4.

Theorem 3.1. *For any integer $j \geq 1$, if f_t is a C^{j+1} family of piecewise expanding C^{j+2} unimodal maps in the topological class of a mixing map f_0 , then there is $\hat{\epsilon} > 0$ so that for any $C^{j+\text{Lip}}$ function ψ the map $R(t) = \int \psi \rho_t dx$ is C^j in $(-\hat{\epsilon}, \hat{\epsilon})$.*

As an immediate corollary of Theorem 3.1 and Proposition A.1, we recover the first claim of [4, Theorem 5.1] if ψ is $C^{1+\text{Lip}}$ (we do not need the assumption $\partial f_t|_{t=0} = X \circ f_0$ used in [4, Theorem 5.1]):

Corollary 3.2. *Assume that f_t is a C^2 family of piecewise expanding C^3 unimodal maps, where f_0 is a good mixing map. If f_t is tangent to the topological class of f_0 then for any $C^{1+\text{Lip}}$ function $\psi : I \rightarrow \mathbb{C}$, the map $R(t) = \int \psi d\mu_t$ is differentiable at $t = 0$.*

Proof of Theorem 3.1. Fix $\psi \in C^{1+\text{Lip}}$, recall the notation h_t from (1), put

$$(14) \quad g_{s,t}(y) = \frac{\exp(s\psi(h_t(y)))}{|f'_t(h_t(y))|}, \quad y \in I \setminus \{c\},$$

and consider the transfer operator

$$(15) \quad \tilde{\mathcal{L}}_{s,t}\varphi(x) = \sum_{f(y)=x} g_{s,t}(y)\varphi(y).$$

Note that $\mathcal{L}_0 = \tilde{\mathcal{L}}_{0,0}$ is the usual transfer operator for f_0 . We let \mathcal{L}_0 act on BV_p for any fixed $p \in [1, p_0)$, where p_0 , depending on f_0 through $\inf |f'_0|$ and $\sup |f'_0|$, is given by the main Theorem of [17], which says that \mathcal{L}_0 on BV_p has spectral radius 1, essential spectral radius < 1 , and 1 is the only eigenvalue of modulus 1. Furthermore, 1 is a simple eigenvalue with an eigenvector ρ_0 which is strictly positive on $[c_2, c_1]$. The corresponding fixed vector ν_0 of \mathcal{L}_0^* is just Lebesgue measure dx . We normalise so that $\int \rho_0 d\nu_0 = 1$ and $\nu_0(I) = 1$. (Of course, $\mu_0 = \rho_0 dx$.) We shall view $\tilde{\mathcal{L}}_{s,t}$ as a perturbation of \mathcal{L}_0 , more precisely we write

$$(16) \quad \tilde{\mathcal{L}}_{s,t}(\varphi) = \mathcal{L}_0 \left(\frac{g_{s,t}}{g_{0,0}} \varphi \right).$$

Since $g_{0,0}^{-1} = |f'_0| \in BV_p$ the bound (13) implies that $\mathcal{P}_{s,t}(\varphi) = \frac{g_{s,t}}{g_{0,0}} \varphi$ is a bounded operator on BV_p . Clearly, $\mathcal{P}_{0,0} = \text{id}$. Since $\|g_{s,t} - g_{0,0}\|_{BV_1} \rightarrow 0$ as $(s,t) \rightarrow (0,0)$, the operators $\tilde{\mathcal{L}}_{s,t}$ on BV_p have a real positive simple maximal eigenvalue with a spectral gap, uniformly in (s,t) close enough to $(0,0)$, by classical perturbation theory [9]. In particular, the operator $\tilde{\mathcal{L}}_t = \tilde{\mathcal{L}}_{0,t}$ on BV_p has a simple eigenvalue at 1, for the fixed point $\tilde{\rho}_t = \rho_t \circ h_t$, where $\mu_t = \rho_t dx$ is the SRB measure of f_t , and the rest of its spectrum lies in a disc of strictly smaller radius. Note that the fixed point of $\tilde{\mathcal{L}}_t^*$ is the measure ν_t defined by

$$(17) \quad \int \varphi dx = \int \varphi \circ h_t d\nu_t.$$

(By definition ν_t is a probability measure and $\int \tilde{\rho}_t d\nu_t = 1$.)

Consider first the case $j = 1$. Lemma 3.3 below implies that the map $s \mapsto \mathcal{P}_{s,t}$ is C^1 from \mathbb{R} to the Banach space of C^1 maps from $\{|t| < \epsilon\}$ to bounded operators on BV_p , and

$$(18) \quad \partial_s \mathcal{P}_{s,t}|_{s=u} = (\psi \circ h_t) \mathcal{P}_{u,t}, \quad \forall u \in \mathbb{R}.$$

Therefore, $s \mapsto \mathcal{L}_{s,t}$ is C^1 from \mathbb{R} to the Banach space of C^1 maps from $\{|t| < \epsilon\}$ to bounded operators on BV_p , and

$$\partial_s \tilde{\mathcal{L}}_{s,t}|_{s=u}(\varphi) = \tilde{\mathcal{L}}_{u,t}((\psi \circ h_t)\varphi), \quad \forall u \in \mathbb{R}.$$

We are thus in a position to apply classical perturbation theory of an isolated simple eigenvalue (see [9, Ch. VII.1.3] for the analytic case, see e.g. [3, Lemma 3.2] for the differentiable setting). It follows on the one hand that, in a neighbourhood of $(0, 0)$, the maximal ² eigenvalue $\lambda_{s,t} > 0$ of $\tilde{\mathcal{L}}_{s,t}$ acting on BV_p is a C^1 function of s to the space of C^1 maps from $\{|t| < \epsilon\}$ to \mathbb{R} . On the other hand, by “tedious but straightforward calculations” and [9, Ch. VII.1.5, Ch. II.2.2], (to quote [13, (5.2)]), we have

$$\partial_s(\log \lambda_{s,t})|_{s=0} = \int \psi \circ h_t \tilde{\rho}_t d\nu_t = \int \psi d\mu_t$$

(use that $\tilde{\rho}_t$ and ν_t are the fixed eigenvectors of $\tilde{\mathcal{L}}_{0,t}$ and its dual). Since $t \mapsto \partial_s(\log \lambda_{s,t})|_{s=0}$ is a C^1 function in a neighbourhood of zero, we have proved Theorem 3.1 in the case $j = 1$. If $j \geq 2$, apply Lemma 3.4 instead of Lemma 3.3. \square

The following result is the key ingredient in the proof of Theorem 3.1, its proof hinges on Proposition 2.3 and [4, Prop. 2.4]:

Lemma 3.3. *Let f_t be a C^2 family of piecewise expanding C^3 unimodal maps in the topological class of f_0 . For any $p > 1$ there exists $\epsilon_p > 0$ so that for any $\psi : I \rightarrow \mathbb{R}$ which is $C^{1+\text{Lip}}$, the map $s \mapsto g_{s,t}$ defined by (14) is C^1 from \mathbb{R} to the Banach space of C^1 maps from $\{|t| < \epsilon_p\}$ to BV_p . In addition, recalling the notation (1),*

$$(19) \quad \partial_s g_{s,t}|_{s=u} = (\psi \circ h_t) g_{u,t}, \quad \forall u \in \mathbb{R}.$$

In fact, s -analyticity holds in Lemma 3.3, but we shall not need this.

Proof of Lemma 3.3. Fix $p > 1$. For every $x \neq c$, all small t , and all $s_1 < s_2$ in \mathbb{R} , there exists $s_3 \in [s_1, s_2]$ so that

$$(20) \quad g_{s_1,t}(x) - g_{s_2,t}(x) - (s_1 - s_2)\psi(h_t(x))g_{s_2,t}(x) = (s_1 - s_2)^2(\psi(h_t(x)))^2 g_{s_3,t}(x).$$

(Just use the Taylor formula for $s \mapsto g_{s,t}(x)$ and the intermediate value theorem.)

So, to prove both differentiability and (19), it suffices to see that the three maps

$$t \mapsto g_{s,t}, \quad t \mapsto (\psi \circ h_t)g_{s,t}, \quad t \mapsto (\psi \circ h_t)^2 g_{s,t},$$

are C^1 from a neighbourhood of 0 to BV_p , uniformly in s in any compact set $K \subset \mathbb{R}$.

In view of this, we first study the maps $t \mapsto h_t(x)$. By [4, Proposition 2.4], there exists $\tilde{\epsilon} > 0$ so that the set of maps $\{t \mapsto h_t(x), x \in I\}$ is bounded in $C^{1+\text{Lip}}([-\tilde{\epsilon}, \tilde{\epsilon}])$. Differentiating with respect to t the equation $h_t \circ f_0 = f_t \circ h_t$, and setting $\alpha_t = \partial_t h_t \circ h_t^{-1}$, we get

$$\alpha_t(f_t(c)) = \partial_t f_t(c), \quad \partial_t f_t(x) = \alpha_t(f_t(x)) - f'_t(x)\alpha_t(x), \quad \forall x \neq c, |x| < \tilde{\epsilon}.$$

Since $\alpha_t(c) = 0$ this implies $J(f_t, \partial_t f_t) = 0$ for $|t| < \tilde{\epsilon}$ (recall (4)), so, for any fixed

$$\beta \in (1/p, 1) \quad (\text{we may and shall assume also that } \beta < 1/\sqrt{p}),$$

Proposition 2.3 gives C and $\epsilon_p > 0$ so that

$$(21) \quad |\alpha_t|_\beta \leq C, \quad \forall |t| < \epsilon_p.$$

Let α_t^η be the η -regularisation (in the variable x) of α_t , that is the convolution $\alpha_t^\eta(x) = \int \alpha_t(y) \kappa_\eta(x - y) dy$ of α_t with a convolution kernel $\kappa_\eta(x) = \eta^{-1} \kappa(x/\eta)$,

²Of course, $\log \lambda_{s,t}$ is the topological pressure of $\log g_{s,t}$.

where the C^∞ function $\kappa : \mathbb{R} \rightarrow \mathbb{R}_+$ is supported in $[-1, 1]$, and $\int \kappa(x) dx = 1$. Note for further use that (21) gives \tilde{C} so that, for all $|t| < \epsilon_p$,

$$(22) \quad \|\alpha_t^\eta\|_{C^{1+1/p}} \leq \frac{\tilde{C}}{\eta^{1+1/p-\beta}}, \quad \|\alpha_t^\eta\|_\beta \leq \tilde{C}, \quad |\alpha_t^\eta - \alpha_t|_{1/p} \leq \tilde{C}\eta^{\beta-1/p}, \quad \forall \eta \in (0, 1).$$

We now consider $t \mapsto g_{s,t}$. For $x \neq c$, we have

$$(23) \quad \partial_t g_{s,t}(x) = e^{s\psi(h_t(x))} \left[\frac{\psi'(h_t(x))\alpha_t(h_t(x))}{|f'_t(h_t(x))|} - \frac{\partial_t(|f'_t(h_t(x))|)}{|f'_t(h_t(x))|^2} \right],$$

where

$$(24) \quad \partial_t(|f'_t(h_t(x))|) = -\text{sgn}(x)(f''_t(h_t(x))\alpha_t(h_t(x)) + \partial_t f'_t(h_t(x))).$$

We claim that the function $x \mapsto \partial_t g_{s,t}(x)$ has bounded $BV_{1/\beta}$ norm, uniformly in $s \in K$ and $|t| < \epsilon_p$. Indeed, decomposing

$$\partial_t g_{s,t} = b_{s,t} \circ h_t,$$

note that each $h_t : I \rightarrow I$ is a homeomorphism leaving both $[-1, c]$ and $[c, 1]$ invariant, while $b_{s,t}$ is β -Hölder on $[-1, c]$ and $(c, 1]$, uniformly in $s \in K$ and $|t| < \epsilon_p$ (because ψ' is C^β , f_t is a C^2 family of C^3 maps³, and α_t is β -Hölder, uniformly in $|t| < \epsilon_p$), and $\sup_{s \in K, |t| < \epsilon_p} |b_{s,t}(c_+) - b_{s,t}(c_-)| < \infty$ (using $\sup_{|t| < \epsilon_p} \|f_t\|_{\mathcal{B}^{2+\beta}} < \infty$).

To conclude, it suffices to prove that our candidate $b_{s,t} \circ h_t \in BV_p$ is really the t -derivative of $g_{s,t}$ (uniformly in s), that is,

$$(25) \quad \lim_{t_2 \rightarrow t_1} \sup_{s \in K} \left\| \frac{g_{s,t_2} - g_{s,t_1}}{t_2 - t_1} - b_{s,t_1} \circ h_{t_1} \right\|_{BV_p} = 0, \quad \forall |t_1| < \epsilon_p,$$

and that this derivative is continuous in t (uniformly in s), that is,

$$(26) \quad \lim_{\tilde{t} \rightarrow t_1} \sup_{t_2 \in [t_1, \tilde{t}]} \sup_{s \in K} \|b_{s,t_2} \circ h_{t_2} - b_{s,t_1} \circ h_{t_1}\|_{BV_p} = 0, \quad \forall |t_1| < \epsilon_p.$$

We first prove (26). Decomposing

$$(27) \quad b_{s,t_2} \circ h_{t_2} - b_{s,t_1} \circ h_{t_1} = (b_{s,t_2} - b_{s,t_1}) \circ h_{t_2} + b_{s,t_1} \circ h_{t_2} - b_{s,t_1} \circ h_{t_1},$$

we focus first on the second term in the right-hand-side of (27). Let $\delta > 0$ be such that f'_t , f''_t and $\partial_t f'_t$ restricted to $[-1, c]$ and $[c, 1]$, respectively, extend to C^1 functions of x on $[-1 - \delta, c + \delta]$ and $[c - \delta, 1 + \delta]$, respectively, for all $|t| < \epsilon_p$. Denote by $b_{s,t}^{\eta,-}$ the function obtained from $b_{s,t}$ by substituting α_t with α_t^η , and also ψ' , and the extensions to $[-1 - \delta, c + \delta]$ of $f''_t|_{[-1, c]}$, $\partial_t f'_t|_{[-1, c]}$, with their x -convolutions with κ_η , for small $\eta > 0$ (to be determined later). Define $b_{s,t}^{\eta,+}$ similarly, using $[c - \delta, 1 + \delta]$, and set $b_{s,t}^\eta(x) = b_{s,t}^{\eta,+}(x)$ if $x > c$ and $= b_{s,t}^{\eta,-}(x)$ if $x < c$. Since $\beta < 1$ and ψ' is Lipschitz, it is easy to see that there exists $\hat{C} > 0$ so that for all $\eta \in (0, 1)$

$$\max\left(\sup_{s \in K} |(b_{s,t_1}^\eta|_{(-\infty, c)})'|_{1/p}, \sup_{s \in K} |(b_{s,t_1}^\eta|_{(c, \infty)})'|_{1/p}\right) \leq \frac{\hat{C}}{\eta^{1+1/p-\beta}}, \quad \forall |t_1| < \epsilon_p.$$

(Use the first two estimates of (22), and the analogous bounds for the regularisations of ψ' and f''_t , $\partial_t f'_t$.) Therefore, by the fundamental theorem of calculus and the

³This implies in particular that $x \mapsto \partial_t f_t$ is C^2 in x , uniformly in t and $\partial_x \partial_t f_t = \partial_t f'_t$.

Hölder (or Jensen) inequality, there exists $\bar{C} > 0$ so that for all $s \in K$, all $|t_1| < \epsilon_p$, $|t_2| < \epsilon_p$, all $\eta \in (0, 1)$, and any $x_0 < x_1 < \dots < x_N \leq c$,

$$\begin{aligned}
& \sum_{i=0}^{N-1} |b_{s,t_1}^\eta(h_{t_2}(x_i)) - b_{s,t_1}^\eta(h_{t_1}(x_i)) - b_{s,t_1}^\eta(h_{t_2}(x_{i+1})) + b_{s,t_1}^\eta(h_{t_1}(x_{i+1}))|^p \\
&= \sum_i \left| \int_{t_1}^{t_2} \partial_t(b_{s,t_1}^\eta(h_t(x_i))) dt - \int_{t_1}^{t_2} \partial_t(b_{s,t_1}^\eta(h_t(x_{i+1}))) dt \right|^p \\
&\leq \sum_i \int_{t_1}^{t_2} |(b_{s,t_1}^\eta)'(h_t(x_i))\alpha_t(h_t(x_i)) - (b_{s,t_1}^\eta)'(h_t(x_{i+1}))\alpha_t(h_t(x_{i+1}))|^p dt \\
&= \int_{t_1}^{t_2} \sum_i |(b_{s,t_1}^\eta)'(h_t(x_i))\alpha_t(h_t(x_i)) - (b_{s,t_1}^\eta)'(h_t(x_{i+1}))\alpha_t(h_t(x_{i+1}))|^p dt \\
&\leq |t_2 - t_1| \left(\sup_{x \in (-\infty, c)} |(b_{s,t_1}^\eta)'(x)| \sup_t \|\alpha_t \circ h_t\|_{BV_p} + \sup_{x,t} |\alpha_t| |(b_{s,t_1}^\eta)'|_{(-\infty, c)}|_{1/p} \right) \\
(28) \quad &\leq \bar{C} \frac{|t_2 - t_1|}{\eta^{1+1/p-\beta}}.
\end{aligned}$$

(We used (21) in the last inequality.) The same bounds hold for $c \leq x_0 < x_1 < \dots < x_N$, and it is easy to estimate the jump of $b_{s,t_1}^\eta \circ h_{t_2} - b_{s,t_1}^\eta \circ h_{t_1}$ at $x = c$ uniformly in s and t_1, t_2 .

We next analyse the contribution of $b_{s,t_1} - b_{s,t_1}^\eta$ to the second term of (27). For this, observe that if h is an orientation preserving homeomorphism fixing c and b is β -Hölder on $[-\infty, c]$ and $[c, \infty]$, then $\|b \circ h\|_{BV_p} \leq |b|_{(-\infty, c)}|_\beta + |b|_{[c, \infty)}|_\beta + |b(c_+) - b(c_-)|$. Then, the last bound of (22) and its analogue for the η -regularisation of ψ' , f_t'' and $\partial_t f_t'$ give a constant C' so that for all $|t_1| < \epsilon_p$, $|t_2| < \epsilon_p$ and $\eta \in (0, 1)$

$$\begin{aligned}
\sup_{s \in K} \|(b_{s,t_1} - b_{s,t_1}^\eta) \circ h_{t_2} - (b_{s,t_1} - b_{s,t_1}^\eta) \circ h_{t_1}\|_{BV_p} &\leq 2 \sup_{s,t} \|(b_{s,t_1} - b_{s,t_1}^\eta) \circ h_t\|_{BV_p} \\
(29) \quad &\leq C' \eta^{\beta-1/p}.
\end{aligned}$$

Taking $\xi \in (0, 1)$ and setting $\eta = (t_2 - t_1)^{\frac{\xi}{1+1/p-\beta}}$, we get from (28-29) that $\lim_{\tilde{t} \rightarrow t_1} \sup_{t_2 \in [t_1, \tilde{t}]} \sup_{s \in K} \|b_{s,t_1} \circ h_{t_2} - b_{s,t_1} \circ h_{t_1}\|_{BV_p} = 0$.

To analyse the first term of (27), we start by noticing that since $t \mapsto \partial_t h_t$ is Lipschitz, there exists a set $\mathcal{D}_p \subset (-\epsilon_p, \epsilon_p)$ of full Lebesgue measure so that $\partial_t h_t$ is differentiable at all t in \mathcal{D}_p . Differentiating twice $f_t \circ h_t(x) = h_t \circ f(x)$ with respect to t and ⁴ setting $\alpha_t^2 = \partial_{tt}^2 h_t \circ h_t^{-1}$, we obtain for all $x \neq c$ and all $t \in \mathcal{D}_p$ that

$$(30) \quad f_t''(x)\alpha_t(x)^2 + 2\partial_t f_t'(x)\alpha_t(x) + \partial_{tt} f_t(x) = \alpha_t^2(f_t(x)) - f_t'(x)\alpha_t^2(x).$$

The left-hand-side of the above TCE is β -Hölder in $[-1, c]$ and $[c, 1]$ and continuous in I , since $\alpha_t(c) = 0$ for every small t , so it is β -Hölder continuous. Therefore, by Proposition 2.3, there exist $\epsilon_p > 0$ and a constant C'' so that

$$(31) \quad |\alpha_t^2|_\beta \leq C'', \quad \forall t \in \mathcal{D}_p.$$

⁴This is similar the proof of [4, Proposition 2.4], but we will make a more careful analysis of what was called F_i there.

The fundamental theorem of calculus holds for the Lipschitz (and therefore almost everywhere differentiable) function $t \mapsto b_{s,t}$ and gives

$$(32) \quad (b_{s,t_2} - b_{s,t_1})h_{t_2}(x) = \int_{t_1}^{t_2} \partial_t b_{s,t}(h_{t_2}(x)) dt, \forall x \neq c.$$

The first term of (27) may then be estimated via the Hölder inequality and the fundamental theorem of calculus (32), as in (28), but exploiting (31) instead of using η -regularisation. Details are left to the reader.

Finally, to show (25), start from

$$g_{s,\tilde{t}}(x) - g_{s,t_1}(x) - (\tilde{t} - t_1)b_{s,t_1}(h_{t_1}(x)) = \int_{t_1}^{\tilde{t}} (b_{s,t_2}(h_{t_2}(x)) - b_{s,t_1}(h_{t_1}(x))) dt_2,$$

for all $x \neq c$, and use the Hölder inequality and (26) (details are left to the reader).

The analysis of the maps $t \mapsto (\psi \circ h_t)g_{s,t}$ and $t \mapsto (\psi \circ h_t)^2 g_{s,t}$ goes along exactly the same lines. \square

For the higher regularity statement in Theorem 3.1, we use the following result (again, analyticity in s holds):

Lemma 3.4. *Let $j \geq 2$. Let f_t be a C^{j+1} family of piecewise expanding C^{j+2} unimodal maps in the topological class of f_0 . For any $p > 1$ there exists $\epsilon_p > 0$ so that for any $\psi : I \rightarrow \mathbb{R}$ which is $C^{j+\text{Lip}}$, the map $s \mapsto g_{s,t}$ defined by (14) is C^1 from \mathbb{R} to the space of C^j maps from $\{|t| < \epsilon_p\}$ to BV_p , and, recalling (1), $\partial_s g_{s,t}|_{s=u} = (\psi \circ h_t)g_{u,t}$.*

Proof. Since the family f_t is C^{j+1} , the set $\{t \mapsto h_t(x), x \in I\}$ is bounded in $C^{j+\text{Lip}}$ by [4, Proposition 2.4]. Let $\beta \in (1/p, 1)$ (with $\beta < 1/\sqrt{p}$, say). Assume first $j = 2$. Then, by (30), the function $\alpha_t^2 = \partial_{tt}^2 h_t \circ h_t^{-1}$, is well-defined for all $|t| < \epsilon_p$ and there exists C so that $|\alpha_t^2|_\beta \leq C$ for every $|t| < \epsilon_p$. For $j \geq 3$, a higher order TCE similar to (30) gives that $\alpha_t^j = \partial_{tj}^j h_t \circ h_t^{-1}$ is β -Hölder for all $|t| < \epsilon_p$. We put $\alpha_t^1 = \alpha_t$.

Then, computing $\partial_{tj}^j g_{s,t}(x)$ at $x \neq c$ gives $b_{s,t}^{(j)}(h_t(x))$, where $b_{s,t}^{(j)}$ is an expression involving derivatives of order at most j of $\psi(x)$, functions α_t^ℓ , for $1 \leq \ell \leq j$, and derivatives (in x, t , or mixed) of total order at most j of $f_t'(x)$, in the numerator, and $|f_t'(x)|^m$ for $m \geq 1$ in the denominator. Our differentiability assumptions on ψ and the family f_t then allow us to proceed as in the proof of Lemma 3.3 (using Taylor series of higher order). \square

4. RECOVERING THE LINEAR RESPONSE FORMULA

Here we give a slightly different proof of the differentiability of $R(t) = \int \psi d\mu_t$, where μ_t is the SRB measure of f_t , still relying heavily on Proposition 2.3 (via Lemma 4.2). The advantage with respect to Theorem 3.1 is that we recover the formula for $\partial_t R(t)|_{t=0}$, and we need only assume that ψ is C^0 . (In particular, this gives a new proof of [4, Theorem 5.1].) We also get new information in Corollary 4.4 by combining Theorems 3.1 and 4.1.

We need notation. By [1, Proposition 3.3], we may decompose the invariant density of a piecewise expanding C^3 unimodal mixing map f_t as $\rho_t = \rho_{\text{reg},t} + \rho_{\text{sal},t}$,

where $\rho_{reg,t} \in BV \cap C^0$, $\rho'_{reg,t} \in BV$, and

$$\rho_{sal,t} = \sum_{k=1}^{M_f} s_{k,t} H_{c_{k,t}}.$$

(Here, $H_u(x)$ denotes the Heaviside function $H_u(x) = -1$ if $x < u$, $H_u(x) = 0$ if $x > u$ and $H_u(u) = -1/2$.) If $M_f = \infty$ then it is not difficult to show that (see e.g. [1, 4], noting that if $c_{1,t}$ is preperiodic but not periodic our notation is slightly different than the notation there)

$$(33) \quad s_{k,t} = \frac{s_{1,t}}{(f^{k-1})'(c_{1,t})}, \quad \forall k \geq 1.$$

We simply write $\rho_0 = \rho = \rho_{reg} + \rho_{sal}$.

To compute the formula for the derivative, we shall assume, as in [4], that $v = \partial_t f_t|_{t=0}$ is of the form $v = X \circ f_0$ for a C^2 function $X : I \rightarrow \mathbb{R}$.⁵

Theorem 4.1. *Let f_t be a C^2 family of piecewise expanding C^3 unimodal maps. Assume that f_0 is good and mixing, that f_t is tangent to the topological class of f_0 , and that $v = \partial_t f_t|_{t=0} = X \circ f_0$ for a C^2 function X . Then, as Radon measures,*

$$(34) \quad \lim_{t \rightarrow 0} \frac{\mu_t - \mu_0}{t} = -\alpha \rho'_{sal} - (\text{id} - \mathcal{L}_0)^{-1} (X' \rho_{sal} + (X \rho_{reg})') dx,$$

where the function α is given by (4), and the operator $\mathcal{L}_0 = \tilde{\mathcal{L}}_{0,0}$ is defined by (15). In addition, α is β -Hölder for any $\beta < 1$.

Proof. Set $f = f_0$ for convenience. By Proposition A.1, we can assume that f_t lies in the topological class of f , denoting the conjugacies by h_t as usual. The transfer operator $\tilde{\mathcal{L}}_t = \tilde{\mathcal{L}}_{0,t}$ for f and the weight $|f'_t \circ h_t|^{-1}$ (recall (15)) is conjugated to the transfer operator for f_t and $|f'_t|^{-1}$ defined by

$$\mathcal{L}_t \varphi(x) = \sum_{f_t(y)=x} \frac{\varphi(y)}{|f'_t(y)|},$$

via

$$(35) \quad \tilde{\mathcal{L}}_t(\varphi \circ h_t) = \mathcal{L}_t(\varphi) \circ h_t.$$

Since h_t is a homeomorphism, it gives rise to an isometry of BV_p , and (35) together with the main Theorem of [17] applied to \mathcal{L}_t imply that there exist $p_0 > 1$ so that for any $p \in [1, p_0)$ there exists $\epsilon_p > 0$ so that for all $|t| < \epsilon_p$ the operator $\tilde{\mathcal{L}}_t$ acting on BV_p has a maximal eigenvalue equal to 1, which is simple, and the rest of the spectrum lies in a disc of strictly smaller radius (i.e., $\tilde{\mathcal{L}}_t$ has a spectral gap). The fixed points of $\tilde{\mathcal{L}}_t$ and its dual, $\tilde{\rho}_t = \rho_t \circ h_t$ and ν_t from (17), were introduced in the proof of Theorem 3.1. We can alternatively see that $\tilde{\mathcal{L}}_t$ has a spectral gap on BV_p by noting that it is a small multiplicative perturbation of $\mathcal{L}_0 = \tilde{\mathcal{L}}_0$ on BV_p : Recall that $\lim_{t \rightarrow 0} \|g_{0,t} - g_{0,0}\|_{BV_1} = 0$ and (16). Observe for further use that this implies $\lim_{t \rightarrow 0} \|\tilde{\rho}_t - \rho_0\|_{BV} = 0$.

From now on, we fix $p \in (1, p_0)$.

We next show that $t \mapsto \tilde{\rho}(t) \in BV_p$ and $t \mapsto \nu_t \in BV_p^*$ are differentiable at $t = 0$. By [4, Prop. 2.4, Cor. 2.6] v is horizontal for f_0 , $t \mapsto h_t(x)$ is differentiable, uniformly in $x \in I$, and $\alpha = \partial_t h_t|_{t=0}$ is continuous, with $\alpha(c) = 0$, $\alpha(c_1) =$

⁵See also the beginning of [16, Section 17].

$X(c)$, and α is the unique bounded solution (4) to the TCE (3). In addition, Proposition 2.3 gives that α is β -Hölder for arbitrary $\beta < 1$ (we shall take $\beta \in (1/p, 1/\sqrt{p})$).

Our assumptions on f_t then imply that v' is C^1 and the following operator is bounded on BV_p :

$$(36) \quad \mathcal{M}\varphi(x) = - \sum_{f(y)=x} \frac{f''(y)\alpha(y) + v'(y)}{|f'(y)|f'(y)} \varphi(y).$$

(Write \mathcal{M} as \mathcal{L}_0 composed with a multiplication operator, like in (16) and use (13).) Lemma 4.2 below easily implies that $t \mapsto \tilde{\mathcal{L}}_t$ is differentiable as an operator on BV_p , and that

$$(37) \quad \partial_t \tilde{\mathcal{L}}_t|_{t=0} = \mathcal{M}.$$

As in the proof of Theorem 3.1, perturbation theory then gives that $t \mapsto \tilde{\rho}_t \in BV_p$ and $t \mapsto \nu_t \in BV_p^*$ are differentiable at $t = 0$.

We next show that $t \mapsto \mu_t = \rho_t dx$ is differentiable as a Radon measure, exploiting the formula for \mathcal{M} to get the claimed formula for $\partial_t \mu_t|_{t=0}$. Fix $\psi : I \rightarrow \mathbb{C}$ continuous. Since $\tilde{\rho}_0 = \rho_0$, we can decompose

$$(38) \quad \int \psi \rho_t dx - \int \psi \rho_0 dx = \int \psi \rho_t dx - \int \psi \tilde{\rho}_t dx + \int \psi \tilde{\rho}_t dx - \int \psi \tilde{\rho}_0 dx.$$

We shall now see that

$$(39) \quad \lim_{t \rightarrow 0} \frac{\int \psi (\rho_t - \rho_t \circ h_t) dx}{t} = - \int \psi \alpha \rho'_0.$$

In view of (39), note first that $s_{k,t} \rightarrow s_k$ as $t \rightarrow 0$: Indeed, since $\lim_{t \rightarrow 0} \|\tilde{\rho}_t - \rho_0\|_{BV} = 0$ we have (in BV)

$$(40) \quad \lim_{t \rightarrow 0} \tilde{\rho}_t = \lim_{t \rightarrow 0} \left(\rho_{reg,t} \circ h_t + \sum_{k=1}^{M_f} s_{k,t} H_{c_k} \right) = \tilde{\rho}_0 = \rho_{reg} + \sum_{k=1}^{M_f} s_k H_{c_k}.$$

(We gave another proof of $\lim_{t \rightarrow 0} s_{k,t} = s_k$ in Step 1 of [4, Proof of Theorem 5.1].)

Decompose $\rho_t - \rho_t \circ h_t$ in (39) into $\rho_{sal,t} - \rho_{sal,t} \circ h_t + \rho_{reg,t} - \rho_{reg,t} \circ h_t$. For the singular term, we have in the sense of Radon measures:

$$(41) \quad \lim_{t \rightarrow 0} \frac{\rho_{sal,t} - \rho_{sal,t} \circ h_t}{t} = - \sum_{k=1}^{M_f} \alpha(c_k) s_k \text{Dirac}_{c_k} = -\alpha \rho'_{sal}.$$

(Just use that $s_{k,t} \rightarrow s_k$ and (33), which implies the $s_{k,t}$ decay exponentially in k uniformly in t .⁶)

We claim that the contribution of the regular term $\rho_{reg,t}$ in the decomposition of (39) is $-\int \psi \alpha \rho'_{reg} dx$. Indeed, recall that $\rho'_{reg,t} \in BV$. In fact, we may decompose

$$(42) \quad \rho_{reg,t} = \rho_{regreg,t} + \rho_{regsal,t},$$

with (see Step 2 of [4, Proof of Theorem 5.1]) $\rho_{regreg,t}$ is C^1 (uniformly in t) and

$$\rho'_{regsal,t} = \sum_{k=1}^{M_f} s'_{k,t} H_{c_{k,t}},$$

⁶Claim (41) was also proved in Step 1 of [4, Proof of Theorem 5.1].

where the $s'_{k,t}$ decay exponentially uniformly in t . In fact (see [4, (69), (71)])⁷

$$s'_{k,t} = E'_{k,t} - E_{k,t}, \text{ with } E'_{k,t} = \frac{s'_{k-1,t}}{(f'_t(c_{k-1,t}))^2}, \quad k \geq 2,$$

and

$$\begin{aligned} E_{k,t} &= \frac{s_{k-1,t} f''_t(c_{k-1,t})}{(f'_t(c_{k-1,t}))^3}, \quad k \geq 2, \quad E'_{1,t} = -\frac{(\rho_{reg,t})'(c)}{(f'_t(c_-))^2} + \frac{(\rho_{reg,t})'(c)}{(f'_t(c_+))^2} \\ E_{1,t} &= \left(-\frac{\rho_{reg,t}(c) f''_t(c_-)}{(f'_t(c_-))^3} + \frac{\rho_{reg,t}(c) f''_t(c_+)}{(f'_t(c_+))^3} \right) \\ &\quad + \sum_{k \geq 2, c_{k-1,t} > c} s_{k-1,t} \left(\frac{f''_t(c_-)}{(f'_t(c_-))^3} - \frac{f''_t(c_+)}{(f'_t(c_+))^3} \right). \end{aligned}$$

The argument above giving $s_{k,t} \rightarrow s_k$ also yields $s'_{k,t} \rightarrow s'_k$ (just differentiate once). Therefore, just like in (41), we have

$$\begin{aligned} (43) \quad \lim_{t \rightarrow 0} \int \psi \frac{\rho_{regsal,t} - \rho_{regsal,t} \circ h_t}{t} dx &= - \sum_{k=1}^{M_f} \int \alpha(x) s'_k \psi(x) H_{c_k}(x) dx \\ &= \int \psi \alpha \rho'_{regsal} dx. \end{aligned}$$

In view of handling the term $\rho_{regreg,t}$ from (42), observe that

$$(44) \quad \lim_{t \rightarrow 0} \|\varphi - \varphi \circ h_t\|_{BV} = 0, \quad \forall \varphi \in C^1.$$

Indeed, for $\delta > 0$ and any partition $x_0 < \dots < x_i < x_{i+1} < \dots < x_n$ let $N \leq n$ be so that $\min(x_N, \inf_t h_t(x_N)) > 1 - \delta$, and since $|h_t(y) - y| = O(t)$ uniformly in y , take t_0 so that $|x_i - h_t(x_i)| < \delta/N$ for all $i \leq N$ and $|t| < t_0$. Then use

$$\begin{aligned} &\sum_{i=0}^{n-1} |\varphi(x_i) - \varphi(x_{i+1}) - \varphi(h_t(x_i)) + \varphi(h_t(x_{i+1}))| \leq 2 \sum_{i=0}^N |\varphi(x_i) - \varphi(h_t(x_i))| \\ &+ \sum_{i=N+1}^{n-1} |\varphi(x_i) - \varphi(x_{i+1})| + \sum_{i=N+1}^{n-1} |\varphi(h_t(x_i)) - \varphi(h_t(x_{i+1}))|. \end{aligned}$$

Since the C^1 norm of $\rho_{regreg,t}$ is bounded uniformly in t , (44) easily implies that

$$(45) \quad \lim_{t \rightarrow 0} \|\rho_{regreg,t} - \rho_{regreg}\|_{BV} = 0.$$

(Note for the record that, since $\|\tilde{\rho}_t - \rho_0\|_{BV} = 0$ and $s_{k,t} \rightarrow s_k$, $s'_{k,t} \rightarrow s'_k$, with t -uniformly exponentially decaying $s_{k,t}$, $s'_{k,t}$, this implies $\lim_{t \rightarrow 0} \|\rho_{reg,t} - \rho_{reg}\|_{BV} = 0$.⁸) Then, by the mean value theorem and the x -uniform differentiability of

⁷If c is periodic then $(\rho_{reg,t})'(c)$ may be undefined, but $(\rho_{reg,t})'(c_{\pm})$ are both defined.

⁸ $\lim_{t \rightarrow 0} \|\rho_{reg,t} - \rho_{reg}\|_{BV} = 0$ can also be proved from Keller-Liverani techniques on the spaces in [4], without using the fact that $\tilde{\rho}_t \rightarrow \rho_0$ in the BV norm. This implies $s'_{1,t} \rightarrow s'_1$, like we proved $s_{1,t} \rightarrow s_1$ in [4, Proof of Theorem 5.1].

$$t \mapsto h_t(x)$$

$$\begin{aligned}
& \lim_{t \rightarrow 0} \int \psi \frac{\rho_{regreg,t} - \rho_{regreg,t} \circ h_t}{t} dx \\
&= \lim_{t \rightarrow 0} \int \psi(x) \frac{\rho_{regreg,t}(x) - \rho_{regreg,t}(h_t(x))}{x - h_t(x)} \frac{x - h_t(x)}{t} dx \\
&= \lim_{t \rightarrow 0} \int \psi(x) \rho'_{regreg,t}(x_t) \frac{x - h_t(x)}{t} dx \\
&= - \lim_{t \rightarrow 0} \int \psi(x) \rho'_{regreg,t}(x_t) \alpha(x) dx \\
(46) \quad &= - \int \psi(x) \rho'_{regreg,0}(x) \alpha(x) dx,
\end{aligned}$$

where x_t is in the interval between x and $h_t(x)$, and we used in the last line that $\rho'_{regreg,t}$ is continuous on the compact interval I , uniformly in t , together with (45). Putting (41-43-46) together, we find (39).

We now turn to the estimation of the term $(\int \psi \tilde{\rho}_t dx - \int \psi \tilde{\rho}_0 dx)/t$ from (38). In view of this, note that (37) implies that (as operators on BV_p)

$$\partial_t(z - \tilde{\mathcal{L}}_t)^{-1}|_{t=0} = (z - \mathcal{L}_0)^{-1} \mathcal{M}(z - \mathcal{L}_0)^{-1}.$$

Therefore, writing the spectral projectors as Cauchy integrals, we get by a simple residue computation, since $(z - \mathcal{L}_0)^{-1} \rho_0 = \rho_0/(z - 1)$,

$$(47) \quad \partial_t(\nu_t(\rho_0) \tilde{\rho}_t)|_{t=0} = (\text{id} - \mathcal{L}_0)^{-1} (\text{id} - \Pi_0) \mathcal{M} \rho_0,$$

where $\Pi_0(\varphi) = \rho_0 \int \varphi dx$.

Next, we claim that we have (in BV_p)

$$(48) \quad -\alpha \rho'_{reg} + (\text{id} - \mathcal{L}_0)^{-1} (\text{id} - \Pi_0) (\mathcal{M} \rho_0) = -(\text{id} - \mathcal{L}_0)^{-1} (\text{id} - \Pi_0) (X' \rho_0 + X \rho'_{reg}).$$

(Recall from [1, Proof of Proposition 4.4] that $\Pi_0(X' \rho_0 + X \rho'_{reg}) = 0$.) Since the TCE (3) implies, using $v' = (X' \circ f) \cdot f'$,

$$\mathcal{M} \rho_0(x) = (X(x) - \alpha(x)) \sum_{f(y)=x} \frac{f''(y)}{|f'(y)| f'(y)^2} \rho_0(y) - X'(x) \rho_0(x),$$

to prove (10), it suffices to show

$$-\alpha \rho'_{reg} + (\text{id} - \mathcal{L}_0)^{-1} (\text{id} - \Pi_0) (X - \alpha) (\widetilde{\mathcal{M}} \rho_0) = -(\text{id} - \mathcal{L}_0)^{-1} (\text{id} - \Pi_0) (X \rho'_{reg}),$$

where $\widetilde{\mathcal{M}} \varphi(x) = \sum_{f(y)=x} \frac{f''(y)}{|f'(y)| f'(y)^2} \varphi(y)$. Note that ⁹ if $x \in [-1, c_1]$ is not along the postcritical orbit we have, using $(\rho_{reg})'(y) = (\rho_0)'(y)$ if y is not on the postcritical orbit,

$$(49) \quad (\rho_{reg})'(x) = (\rho_0)'(x) = (\mathcal{L}_0(\rho_0))'(x) = \sum_{f(y)=x} \frac{(\rho_{reg})'(y)}{|f'(y)| f'(y)} - \frac{\rho_0(y) f''(y)}{|f'(y)| (f'(y))^2}.$$

Therefore, for any $x \in I$ which is not on the postcritical orbit

$$\widetilde{\mathcal{M}}(\rho_0)(x) = \sum_{f(y)=x} \frac{\rho'_{reg}(y)}{|f'(y)| f'(y)} - \rho'_{reg}(x).$$

⁹This is [4, (70)].

In other words, we have (in BV_p)

$$\widetilde{\mathcal{M}}(\rho_0) = \bar{\mathcal{M}}(\rho'_{reg}) - \rho'_{reg},$$

where $\bar{\mathcal{M}}\varphi(x) = \sum_{f(y)=x} \frac{\varphi(y)}{|f'(y)|f'(y)}$. So we have reduced the claim (48) to

$$-\alpha\rho'_{reg} - (\text{id} - \mathcal{L}_0)^{-1}(\text{id} - \Pi_0)(\mathcal{L}_0 - \text{id})(\alpha\rho'_{reg}) = 0,$$

that is, using $\Pi_0\mathcal{L}_0 = \Pi_0$,

$$(\mathcal{L}_0 - \text{id})(\alpha\rho'_{reg}) = (\text{id} - \Pi_0)(\mathcal{L}_0 - \text{id})(\alpha\rho'_{reg}) = (\mathcal{L}_0 - \text{id})(\alpha\rho'_{reg}).$$

Finally, since $\rho_0 \in BV_p$, and since $t \mapsto \nu_t$ and $t \mapsto \tilde{\rho}_t$ are differentiable in BV_p and BV_p^* , respectively, we have (in BV_p)

$$(50) \quad \partial_t(\nu_t(\rho_0)\tilde{\rho}_t)|_{t=0} = \partial_t(\nu_t(\rho_0))|_{t=0}\rho_0 + \partial_t(\tilde{\rho}_t)|_{t=0}.$$

Take the Lebesgue average of both sides of (50). Since $\partial_t \int \tilde{\rho}_t dx = 0$ (because each $\tilde{\rho}_t dt$ is a probability), and since

$$- \int (\text{id} - \mathcal{L}_0)^{-1}(X'\rho_{sal} + (X\rho_{reg})') dx = 0$$

(use again $\Pi_0(X'\rho_0 + X\rho'_{reg}) = 0$), we find that $\partial_t(\nu_t(\rho_0))|_{t=0} \int \rho_0 dx = 0$. Therefore, $\partial_t(\nu_t(\rho_0))|_{t=0}$, and putting together (38), (39), (47), (48), and (50), we have proved the theorem. \square

We have (a simplification of Lemma 3.3):

Lemma 4.2. *Let f_t be a C^2 family of piecewise expanding C^3 unimodal maps in the topological class of f_0 . Set $v = \partial_t f_t|_{t=0}$. For any $p > 1$ the map $t \mapsto g_t = \frac{1}{|f'_t \circ h_t|} \in BV_p$ is C^1 in a neighbourhood of 0, and $\partial_t g_t|_{t=0} = -\frac{f_0''\alpha + v'}{|f_0|f_0'}$.*

Proof. Differentiability follows from Lemma 3.3 applied to $\psi \equiv 0$. The value of the derivative is given by (23–24) in the proof of that lemma, since $\partial_t f'_t|_{t=0} = v'$. \square

Remark 4.3. We have the following strengthening of Lemma 4.2 if f_t is a C^3 family of piecewise expanding C^4 unimodal maps: For any $p > 1$ the map $t \mapsto g_t = \frac{1}{|f'_t \circ h_t|} \in BV_p$ is C^2 in a neighbourhood of 0, and $|g_t - g_0 + t \frac{f_0''\alpha + v'}{|f_0|f_0'}| = O(t)$. Recalling (36–37), this implies that $\|\frac{\tilde{\mathcal{L}}_t - \mathcal{L}_0}{t} - \mathcal{M}\|_{BV_p} = O(t)$.

We shall get the following (new) result as a corollary of Theorems 3.1 and 4.1:

Corollary 4.4. *If f_t is a C^2 family of piecewise expanding C^3 unimodal maps in the topological class of f_0 , and if $\partial_t f_t|_{t=0} = X \circ f_0$ for a C^2 function X , then there exists $\epsilon > 0$ so that $t \mapsto \mu_t$ is C^1 from $(-\epsilon, \epsilon)$ to Radon measures.*

In particular, under the assumptions of Corollary 4.4, the Radon measure

$$-\alpha_t \rho'_{sal,t} - (\text{id} - \mathcal{L}_t)^{-1}(X'_t \rho_{sal,t} + (X_t \rho_{reg,t})') dx$$

(recall (34)) is continuous as a function of t . (Here, α_t solves (3) for f_t and $v_t = \partial_s f_s|_{s=t}$, and $X_t \circ f_t = v_t$.) This fact is not clear a priori from the formula.

Remark 4.5. We expect that a careful analysis of the term (39) for C^1 functions ψ would allow to bypass the reference to Theorem 3.1 in the proof of Corollary 4.4.

Proof of Corollary 4.4. We want to show that $t \mapsto \partial_u \mu_u|_{u=t} = \tilde{\mu}_t$ is continuous: We know that $\tilde{\mu}_t$ exists for all small t (as a Radon measure) by Theorem 4.1. Clearly, $|\int \psi d\tilde{\mu}_t| \leq C \sup |\psi|$ for all continuous ψ and all small enough t .

Assume for a contradiction that $t \mapsto \tilde{\mu}_t$ is discontinuous at t_0 . This means that there exist $\psi \in C^0$, with $\sup |\psi| = 1$, $\delta > 0$, and a sequence t_m with $|t_m - t_0| < 1/m$, so that $|\int \psi d\tilde{\mu}_{t_0} - \int \psi d\tilde{\mu}_{t_m}| > \delta$ for all m . Take $\tilde{\psi} \in C^1$ so that $\sup |\psi - \tilde{\psi}| < \delta/4$. Then $|\int \tilde{\psi} d\tilde{\mu}_{t_0} - \int \tilde{\psi} d\tilde{\mu}_{t_m}| > \delta/2$ for all m . But Theorem 3.1 implies $|\int \tilde{\psi} d\tilde{\mu}_{t_0} - \int \tilde{\psi} d\tilde{\mu}_{t_m}| < \delta$ if m is large enough, a contradiction. \square

APPENDIX A. A CONSEQUENCE OF THE KELLER-LIVERANI BOUNDS FROM [4]

We state here for the record an immediate corollary of [4, Proposition 3.3] which was based on results in [12] (see Remark 2.1 and note that the assumptions below imply $\sup_I |f_t - \tilde{f}_t| = O(t^2)$):

Proposition A.1. *Let f_t be a C^2 family of piecewise expanding C^2 unimodal maps. Assume that f_0 is mixing and good, and that f_t is tangent to the topological class of f_0 , denoting by \tilde{f}_t a family in the topological class of f_0 with $\tilde{f}_0 = f_0$ and $\partial f_t|_{t=0} = \partial \tilde{f}_t|_{t=0}$.*

Let $\mu_t = \rho_t dx$ and $\tilde{\mu}_t = \tilde{\rho}_t dx$ be the SRB measures of f_t and \tilde{f}_t , respectively. Then for any $\xi < 2$ there exists $C > 0$ so that for all small t

$$\|\rho_t - \tilde{\rho}_t\|_{L^1(Leb)} \leq C|t|^\xi.$$

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